# **Optimal Control for High-Precision Atom Interferometry**

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## **1** Atomic Fountain Interferometer



10 m atomic fountain at Stanford: launch BEC cloud of <sup>87</sup>Rb atoms.



laser pulses implement atomic beamsplitter and mirror

#### Applications

- Gravitational detection:  $\Delta \phi = ma \Delta z T/\hbar$ ;
- gravitational wave detector, satellite-based sensing of underground structures • test of equivalence principle

# **2** Pulse Schemes

### Hamiltonian couples momentum and internal state: conservation of angular momentum



- Raman: internal state of atom changes [2] adiabatic elimination of  $|i\rangle$ ; or rapid adiabatic passage
- Bragg: internal state of atom unchanged generalization: different Bragg orders by tuning frequencies [3] adiabatic elimination of  $|i\rangle$ ; or rapid adiabatic passage
- also: atomic lattice waveguides [4]

### Challenges

- require short pulses to avoid Doppler sensitivity
- sequential transitions: accumulating errors
- dephasing from lase intensity dependent Stark shift
- intensity variations across ensemble

## References

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 $\Rightarrow$  optimal control

# **3 Optimal Control Methods**

#### goal: find controls that minimize functional for reaching the target state

$$J_{\mathcal{T}} = 1 - \mathfrak{Re} \left\langle \Psi(\mathcal{T}) \mid \Psi^{\mathsf{tgt}} \right\rangle; \quad |\Psi(\mathcal{T})\rangle = \exp \left| - \Psi(\mathcal{T}) \right\rangle$$

iterative procedure:

- start from guess pulse
- simulate dynamics; calculate pulse update  $\Delta u(t)$
- updated ( "optimized") pulse is guess for next iteration

### Gradient-free: Nelder-Mead Simplex

- systematically vary control parameters to "roll down the landscape"
- numerically cheap: only need to evaluate  $J_T$
- Only works for a small number of free control parameters!

Gradient-based: Krotov's method [6]

auxiliary functional  $J = J_T + \sum_i \frac{\lambda_i}{S_i(t)} \int_0^T (u_i^{(1)}(t) - \underbrace{u_i^{(0)}(t)}_{u_i^{\text{ref}}})^2 dt$ 

update  $\Delta u_i(t) = \frac{S_i(t)}{\lambda_i} \Im \left\{ \chi^{(0)}(t) \left| \frac{\partial H}{\partial u_i(t)} \right| \Psi^{(1)}(t) \right\};$ boundary cond.  $\left| \chi^{(0)}(T) \right\rangle = \frac{\partial J_T}{\partial \langle \Psi^{\text{tgt}} |}$ 

alternative method: gradient ascent (GRAPE) [7]:  $\Delta u_i(t) \propto \frac{\partial J_T}{\partial u_i(t)}$ 

hybrid approach for best results: pre-optimize with Nelder-Mead first, then continue with gradient search.

# **4 Robustness through Ensemble Optimization**

#### idea: sample noise-realizations of Hamiltonian

example: Rydberg gate [8]

 $\Delta_2$  ----- $\Omega_R(t)$  $\Delta_1$  ---- $\Omega_B(t)$ 

dominant noise sources:

- fluctuation of Rydberg level (stray magnetic fields)
- fluctuation in pulse amplitude  $\Rightarrow$  dipole value

compared to best analytical scheme: optimal control reduces gate duration from 800 ns to 100 ns, and is order of magnitude more robust.





$$\hat{H}(t) = \begin{cases} \ddots & \ddots & 0 \\ \ddots & E_{-1}(t)/\hbar & -\Omega_{-n}(t) \\ 0 & -\Omega_{-n}^{*}(t) & E_{0}(t)/\hbar & -\Omega_{-n}^{*}(t) \\ 0 & -\Omega_{+n}^{*}(t) & E_{+n}^{*}(t) \\ 0 & 0 \end{cases}$$



# **6** Outlook

• Use full model with additional levels; no adiabatic elimination

• Ensemble optimization to address *challenges* (see (2) )

Modeling: QNET computer algebra system

Simulation and optimization:

